isotope effects and suicide inactivation in P 450 -mediated epoxidations ${ }^{11}$ and primary kinetic isotope effects in methyl group hydroxylations. ${ }^{12}$

It is important to note that the simple oxene model's apparently wide applicability suggests only that the active oxygen species of cytochrome P 450 oxidations has radical character. It should not be construed as implying that the enzymatic oxygen is a free atomic oxygen. What has been claimed is that, thus far, the reactions of triplet oxene with various organic compounds seem to have a great deal in common with P 450 -mediated oxidations of the same or similar compounds and have provided a qualitative framework within which a consistent set of mechanisms has been built up. The actual relationship of this oxygen with both the heme
iron and the substrate during oxidation remains to be elucidated.
Acknowledgment, Computations were performed on the CDC 7600 computer at Lawrence Berkeley Laboratory and the Computer Resources VAX 11/780 at SRI International. We thank Dr. Dale Spangler for helpful comments and advice and Dr. Tetsuro Oie for performing the ab initio calculations. Support from NIH Grant No. GM 27943-02 and NCI Contract No N01-CP-15730 is gratefully acknowledged. B.A.M. is a Pharmacology Research Associate, National Institute of General Medical Sciences.

Registry No, $\mathrm{CCl}_{4}, 56-23-5 ; \mathrm{CHCl}_{3}, 67-66-3$; cytochrome P 450 9035-51-2.

# Ab Initio Studies of Molecular Geometries, 27. Optimized Molecular Structures and Conformational Analysis of $N^{\alpha}$-Acetyl- $N$-methylalaninamide and Comparison with Peptide Crystal Data and Empirical Calculations 

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#### Abstract

The molecular geometries (bond distances and bond angles) of $N^{\alpha}$-acetyl- $N$-methylalaninamide were refined in seven characteristic areas of its conformational space by ab initio gradient relaxation at the $4-21 \mathrm{G}$ level. The optimized conformations in order of increasing energy are $C_{7}{ }^{\text {eq }}$ (I), $C_{5}$ (II), $C_{7}{ }^{\text {ax }}$ (III), $\beta_{2}$ (IV), $\alpha_{\mathrm{R}}$ (V), $\alpha_{\mathrm{L}}$ (VI), and $\alpha^{\prime}$ (VII). The variations in local geometry found between the conformations investigated are discussed in detail. It is found that comparable bond distances and bond angles in different conformations can vary by $0.025 \AA$ and up to $7.5^{\circ}$, respectively. Small deviations from planarity of the peptide group (up to $8^{\circ}$ ) are found for some of the conformations. The calculations confirm a previously suggested correlation between the $\phi$ and $\psi$ angles in the helical forms of the dipeptide, which is in agreement with observed protein structure data and high-resolution crystal structures of small polypeptides. Details in the refined local geometries confirm that it is reasonable to rationalize this correlation in terms of a dipeptide specific intramolecular interaction between atoms N7 and H18 (Figure 1). This interaction, which is directed perpendicularly to the peptide bond, may be an important contribution to the formation and stability of bend type structures in proteins. The significance of the variations in local geometry for empirical conformational analysis of proteins is discussed.


In the last decade it has become common practice to investigate the conformational and structural features of peptides by computational methods. Model dipeptides and other systems have been studied by limited basis set ab initio methods, ${ }^{1-4}$ as well as by many different semiempirical and empirical methods. ${ }^{5-8}$ These calculations have given us a general understanding of conformational parameters at low-energy minima in the conformational space. However, little detailed information is available for comparing the effects of changes in conformation on the local geometry of the peptide. This is true even when considering high-resolution experimental data found from crystal structures, where a variety of conformations for equivalent molecules are simply not found.

For molecules larger than dipeptides, one must move away from rigorous computational studies into the area of empirical energy calculations, where approximations of forces and geometries are only as good as the data and equations used to define the force field. Empirical methods are necessary if one wishes to study larger molecular systems simply because the rigorous methods are incapable of handling a large number of atoms at a reasonable computational expense. Thus, empirical methods have been used

[^0]to find low-energy conformations of polypeptides containing up to $\sim 14$ amino acids ${ }^{9,10}$ and have been used in protein structure refinement studies. ${ }^{11}$
The study described here on the molecule $N^{\alpha}$-acetyl- $N$ methylalaninamide (Ala) is an attempt to fill in some important gaps in our understanding of peptide structure. One point of interest, for example, concerns the variations in local geometry

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Figure 1, Atom numbering for $N^{\alpha}$-acetyl- $N$-methylalaninamide.
that are encountered when bond lengths and angles of the same type are compared in different energy-optimized conformations. Comparing the geometry-refined extended $(\beta)$ or $C_{5}$ region conformer with the refined $\alpha$-helix ( $\alpha_{\mathrm{R}}$ or $\alpha_{\mathrm{L}}$ ) conformers or the two hydrogen-bonded ring conformers $C_{7}{ }^{e q}$ and $C_{7}{ }^{\text {ax }}$ with the other regions will give us consistent structural trends which can be used to characterize any given conformation of minimum energy. Data of this kind are unavailable by any other method of study.

An additional goal of this paper will be to examine the procedures used in empirical energy calculations of peptides by comparing the $a b$ initio structural and conformational results with empirical predictions, so that obvious deficiencies in the empirical studies can be focused upon and ultimately corrected. Modifications for the improvement of empirical potentials and geometry will be presented elsewhere. ${ }^{12}$

## Computational Procedures

The ab initio calculations reported here were executed by using Paulay's ${ }^{13}$ gradient method, Pulay's ${ }^{14}$ program with the $4-21 \mathrm{G}$ basis set, ${ }^{15}$ and a normal coordinate force relaxation procedure described previously. ${ }^{16}$ The initial starting conformations (i.e., torsional angles) were taken from results of empirical energy calculations (ECEPP). ${ }^{17}$ Additional details characterizing the computational procedures can be found elsewhere. ${ }^{18}$ Estimates for uncertainties, in bond distances and bond angles based on comparisons of calculated with experimental structures, are presented in ref 1 (Table I).

On the basis of the experiences made during this refinement and that of the homologuous glycine compound, ${ }^{1}$ one can conclude that computations of this kind, on relatively complex and floppy molecules, cannot be carried out with similar expectations as equivalent calculations for small, rigid molecules. This is so, because the potential energy surfaces of such systems are relatively flat in many areas of conformational space. Thus significant variations in geometry may be accompanied by nearly negligible changes in energy. For example, a change of $3^{\circ}$ in the $\psi$ angle of conformer I, from $73^{\circ}$ (Tables I and II) to $70^{\circ}$, with simultaneous relaxation of bond distances and bond angles at the new point in conformational space, has been found to correspond to a change in energy of only $0.06 \mathrm{kcal} / \mathrm{mol}$. Similarly, a change of $\phi$ and $\psi$ in II from $-165^{\circ}$ and $167^{\circ}$, respectively, to $-155^{\circ}$ and $157^{\circ}$, was accompanied by an energy change of only $0.5 \mathrm{kcal} / \mathrm{mol}$. When a molecular system is this flexible, the concept of a local conformational energy minimum becomes less meaningful for the thermally averaged, real molecule. Rather, one should think in terms of characteristic areas of the conformational energy surface, in which the thermally averaged system is characterized by large amplitude motions.

In accordance with these considerations, the geometries presented in Table II are not minimum energy conformations in the traditional (i.e., small molecule) sense. Rather, the geometries given in this paper are structures (bond distances and bond angles) relaxed at specific points of the potential energy surface of Ala defined by their torsional angles. These points represent typical conformationally allowed areas in the $\phi-\psi$ space of the alanine dipeptide and were arrived at in the following way. In all the conformations refined, (for details see below) the optimized torsional angles obtained by empirical energy calculations ${ }^{17}$ were used as starting values. For these structures the bond distances and angles were allowed to relax without any geometrical constraints. Torsional angles were also optimized, but in small steps by applying a large value to the corresponding second-order force constant in the stepping procedure. This made it possible for the torsional angles to move away from

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Figure 2, Conformer I of $N^{\alpha}$-acetyl $N$-methylalaninamide.


Figure 3, Conformer II of $N^{\alpha}$-acetyl- $N$-methylalaninamide.


Figure 4. Conformer III of $N^{\alpha}$-acetyl- $N$-methylalaninamide.
their empirical values into the corresponding flat areas of potential energy in $4-21 \mathrm{G}$ space, but it did not necessarily allow them to seek the exact positions of relatively shallow minima. This procedure was followed because the shallow energy surfaces, as considered above, did not at present justify the huge computing expenses needed to locate exact minima and because it is the primary purpose of this paper to compare trends in local geometries (differences in bond distances and bond angles) for different characteristic conformations of Ala such as the $\alpha$-helical, $C_{7}$, or $C_{5}$ conformers.

This then is the exact import of the geometries reported here. The primary structures are highly relaxed at each point of the potential energy surface considered. That is, estimates for the largest deviations of bond distances and bond angles from their true $4-21 \mathrm{G}$ minima at a given combination of torsional angles are approximately $4 \times 10^{-4} \AA$ and $0.5^{\circ}$, respectively. (As a rule, most parameters are relaxed much better than that.) The calculated energies are practically the exact 4-21G values for the given combinations of torsional angles. These energies characterize the specific conformations of Ala considered but, for the reasons specified above, not necessarily at their exact shallow bottom of potential energy. These values are, therefore, approximate but, at the same time, they are immensely more accurate than ab initio energies obtained in the most frequently practiced way, i.e., from totally unrelaxed, "standard" geometries. Uncertainties in the torsional angles, at which the refinement was terminated (i.e., their relations to true $4-21 \mathrm{G}$ minima), are impossible to estimate with confidence because of the reasons discussed above. A very extensive grid and path search, between the various characteristic areas of potential energy, including complete geometry optimization at each point, would be necessary to determine the exact topography of the Ala potential energy. This may be worthwile to do in the future, but it is not the purpose of this paper. In spite of these uncertainties in the


Figure 5, Conformer IV of $N^{\alpha}$-acetyl- $N$-methylalaninamide.


Figure 6, Conformer V of $N^{\alpha}$-acetyl- $N$-methylalaninamide.


Figure 7, Conformer VI of $N^{\alpha}$-acetyl- $N$-methylalaninamide.
torsional angles, the $\phi$ and $\psi$ values given will nevertheless allow for some interesting and definite conclusions (see below), because of the directions in which the 4-21G optimization moved away from the ECEPP conformational energy minima used as the starting positions. Finally, uncertainty estimates for out-of-plane bends are given at approximately $0.5^{\circ}$.

## Results and Discussion

The results of the calculations are presented in Tables I and II. The atom numbering for $N^{\alpha}$-acetyl- $N$-methylalaninamide (Ala) is given in Figure 1. The seven conformations (I-VII), which were geometry optimized for this study, are denoted as $C_{7}{ }^{\text {eq }}(\mathrm{I})$, $C_{5}$ (II), $C_{7}{ }^{\text {ax }}$ (III), $\beta_{2}$ (IV), $\alpha_{\mathrm{R}}$ (V), $\alpha_{\mathrm{L}}$ (VI), and $\alpha^{\prime}$ (VII) and are presented in Figures 2-8, respectively. Optimized Cartesian coordinates of I-VII, total energies, calculated dipole moments, and largest Cartesian residual force components on any atom are listed in Table I. Some selected internal coordinates (bond lengths, bond angles, and dihedral angles) can be found in Table II. In Table III we present a comparison of the $\phi$ and $\psi 4-21 \mathrm{G}$ torsional angles reported here and the corresponding ECEPP (empirical conformational energy program for Peptides ${ }^{19}$ ) minimum energy
(19) Zimmerman, S. S.; Pottle, M. S.: Nemethy, G.; Scheraga, H. A. Macromolecules 1977, 10, 1-9.


Figure 8, Conformer VII of $N^{\alpha}$-acetyl- $N$-methylalaninamide.


Figure 9, A $\phi\left(\mathrm{N}-\mathrm{C}^{\alpha}\right), \psi\left(\mathrm{C}^{\alpha}-\mathrm{C}^{\prime}\right)$ conformational isoenergetic contour diagram for the molecule $N^{\alpha}$-acetyl- $N$-methyl-L-alaninamide. The energy is in kilocalories per mole and the solid dots are energy minima found by using ECEPP, while the circles are the local geometry positions corresponding to the torsional angles obtained by the present $4-21 \mathrm{G}$ analysis. (See text for details concerning the nature of the torsional angles.)
values for Ala and Gly. The ab initio results for Gly were taken from ref 1 , and similar considerations apply to the reported values of its torsional angles as specified for Ala above. To further characterize the displacement of $\phi$ and $\psi$ during the $4-21 \mathrm{G}$ refinement from their starting (ECEPP) values, $\phi$ and $\psi$ angles of Table III are plotted in Figure 9 on an isoenergetic $\phi-\psi$ map for Ala obtained by using ECEPP.
The existence of one or more local conformational minima, not included in the seven structures studied here, cannot be excluded on the basis of this study. However, a more rigorous search of the potential energy surface of this compound is not possible at this time due to computational limitations imposed by the large size of the molecule. In the next sections, details of the refined geometries, and their interpretation with respect to allowed minimum energy states, will be discussed.
$C_{7}{ }^{\text {aq }}$ and $C_{7}{ }^{\text {ax }}$ States, The conformation of lowest energy found by the $4-21 \mathrm{G}$ calculations is the $C_{7}{ }^{\text {eq }}$ conformer (I). ECEPP calculations also find this conformation to be that of lowest energy. ${ }^{19}$ However, compared to $C_{7}{ }^{\text {eq }}, C_{7}{ }_{7}$ ax is much more unstable in the ECEPP calculations ( $8.8 \mathrm{kcal} / \mathrm{mol}$, see Table III) than in $4-21 \mathrm{G}$ space ( $2.6 \mathrm{kcal} / \mathrm{mol}$ ). One reason for this difference in the two results may be found in the fact that rigid geometries were used in the ECEPP calculations, whereas significant changes in backbone parameters resulted from the 4-21G geometry optimization. For example (see Table II), the C6-N7-C9 angle is

Table I. 4-21G Optimized Cartesian Coordinates ( $\AA$ ), Total Energies, $E$ ( $k$ cal/mol), Dipole Moments, $D$ (debye), and largest Residual Cartesian Force Component on Any Atom, $F$ (mdyn), for the Conformations $1-\mathrm{V} 11$ of $N^{\alpha}$-Acetyl- $N$-methylalaninamide ${ }^{\alpha}$


[^3]Table II. Some Selected Internal Coordinates ${ }^{a}$ for the $4-21 \mathrm{G}$ Optimized Structures $1-\mathrm{V} 11$ of $N^{\alpha}$-Acetyl- $N$-methylalaninamide

|  | I | 11 | 11 I | 1 V | V | V1 | VII |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| C4-H1 | 1.077 | 1.077 | 1.077 | 1.077 | 1.077 | 1.077 | 1.077 |
| $\mathrm{C} 4-\mathrm{H} 2$ | 1.082 | 1.083 | 1.083 | 1.083 | 1.083 | 1.083 | 1.083 |
| C4-H3 | 1.083 | 1.083 | 1.082 | 1.083 | 1.083 | 1.083 | 1.083 |
| C11-H14 | 1.080 | 1.083 | 1.082 | 1.079 | 1.079 | 1.079 | 1.082 |
| C11-H15 | 1.081 | 1.079 | 1.082 | 1.081 | 1.082 | 1.082 | 1.079 |
| C11-H16 | 1.081 | 1.081 | 1.077 | 1.082 | 1.084 | 1.078 | 1.082 |
| N7-H8 | 0.995 | 0.997 | 0.995 | 0.996 | 0.995 | 0.994 | 0.994 |
| N17-H18 | 0.998 | 0.993 | 0.998 | 0.992 | 0.991 | 0.991 | 0.993 |
| C6=05 | 1.230 | 1.226 | 1.231 | 1.220 | 1.218 | 1.219 | 1.222 |
| C12=O13 | 1.225 | 1.227 | 1.226 | 1.223 | 1.226 | 1.220 | 1.220 |
| C6-C4 | 1.517 | 1.518 | 1.519 | 1.518 | 1.518 | 1.518 | 1.518 |
| C6-N7 | 1.353 | 1.353 | 1.353 | 1.369 | 1.371 | 1.369 | 1.362 |
| C9-N7 | 1.473 | 1.453 | 1.478 | 1.467 | 1.463 | 1.470 | 1.471 |
| C11-C9 | 1.528 | 1.545 | 1.540 | 1.532 | 1.539 | 1.533 | 1.541 |
| C12-C9 | 1.539 | 1.527 | 1.538 | 1.534 | 1.534 | 1.538 | 1.526 |
| C12-N17 | 1.346 | 1.348 | 1.343 | 1.349 | 1.344 | 1.351 | 1.360 |
| C19-N17 | 1.462 | 1.466 | 1.462 | 1.463 | 1.463 | 1.463 | 1.466 |
| C9-H10 | 1.079 | 1.081 | 1.077 | 1.081 | 1.079 | 1.083 | 1.077 |
| N7-C6-C4 | 115.2 | 114.7 | 114.6 | 114.4 | 114.5 | 114.7 | 114.8 |
| N7-C6-O5 | 122.0 | 122.1 | 123.5 | 122.4 | 121.1 | 122.1 | 121.9 |
| C9-N7-C6 | 121.9 | 121.3 | 125.9 | 121.4 | 122.2 | 121.4 | 120.3 |
| C12-C9-N7 | 109.5 | 106.4 | 112.7 | 110.6 | 114.0 | 111.1 | 108.5 |
| C12-C9-C11 | 110.7 | 110.5 | 111.8 | 110.7 | 108.3 | 110.6 | 111.6 |
| O13-C12-C9 | 122.1 | 121.7 | 120.3 | 121.7 | 119.4 | 121.6 | 122.7 |
| H14-C11-C9 | 109.9 | 111.3 | 109.4 | 109.9 | 108.8 | 109.6 | 110.2 |
| H15-C11-C9 | 110.1 | 108.2 | 109.5 | 109.6 | 110.3 | 109.8 | 108.5 |
| H16-C11-C9 | 110.1 | 109.7 | 110.9 | 110.5 | 110.6 | 109.5 | 111.1 |
| N17-C12-C9 | 114.2 | 115.7 | 116.1 | 115.5 | 117.2 | 115.2 | 115.2 |
| C19-N17-C12 | 120.2 | 120.3 | 119.9 | 120.3 | 120.3 | 120.1 | 119.8 |
| C9-N7-H8 | 117.8 | 115.5 | 116.1 | 117.8 | 118.1 | 118.1 | 118.8 |
| C12-C9-H10 | 108.6 | 109.8 | 104.2 | 107.8 | 106.0 | 106.0 | 106.7 |
| H8-N7-C6-C4 | 4.6 | 0.7 | -5.5 | -6.9 | -7.1 | 8.1 | 5.3 |
| H8-N7-C6-O5 | -175.3 | -179.2 | 174.1 | 173.5 | 173.0 | -173.1 | -175.1 |
| C9-N7-C6-C4 | -177.0 | 178.7 | 176.0 | -172.4 | -173.1 | 174.5 | 176.6 |
| C9-N7-C6-O5 | 3.1 | $-1.2$ | -4.4 | 8.1 | 7.0 | -6.7 | -3.8 |
| C11-C9-N7-C6 | 153.5 | 73.7 | -52.9 | 101.6 | 145.7 | $-64.4$ | 74.9 |
| N17-C12-C9-H10 | -43.2 | 49.2 | -176.5 | -76.7 | -124.1 | -73.8 | -172.4 |
| C12-C9-N7-C6 | -84.6 | -165.7 | 74.6 | -134.2 | -91.9 | 60.8 | -161.7 |
| H14-C11-C9-C12 | 62.0 | 61.9 | 51.4 | 53.9 | 52.7 | 55.6 | 57.4 |
| H15-C11-C9-C12 | -177.8 | -178.3 | 170.4 | 174.4 | 172.9 | 176.0 | 176.8 |
| H16-C11-C9-C12 | $-56.3$ | -58.9 | -69.4 | -65.2 | -66.4 | -63.5 | -63.4 |
| O13-C12-C9-N7 | -106.5 | -14.0 | 119.6 | -144.1 | 176.7 | -144.7 | 124.8 |
| O13-C12-C9-C11 | 15.0 | 107.3 | -112.5 | -19.1 | --59.7 | -18.0 | -111.8 |
| N17-C12-C9-N7 | 73.0 | 167.3 | -62.0 | 38.1 | -5.5 | 40.6 | -55.4 |
| $\mathrm{N} 17-\mathrm{C} 12-\mathrm{C} 9-\mathrm{Cl1}$ | -165.4 | -71.4 | 65.9 | 163.1 | 118.1 | 167.3 | 68.1 |
| H18-N17-C12-C9 | -5.1 | --3.4 | 5.5 | -3.6 | 6.3 | -10.3 | -9.2 |
| H18-N17-C12-O13 | 174.4 | 177.9 | -176.2 | 178.6 | -176.0 | 175.1 | 170.7 |
| $\mathrm{C} 19-\mathrm{N} 17-\mathrm{C} 12-\mathrm{C} 9$ | -179.0 | 178.1 | -178.9 | 176.9 | 179.0 | 178.3 | -178.9 |
| C19-N17-C12-O13 | 0.5 | -0.7 | -0.6 | -0.9 | -3.3 | 3.7 | 1.0 |

${ }^{a}$ Bond distances in angstroms and angles in degrees.
Table III, A Comparison of ECEPP Conformational Positions of Minimum Energy for $N^{\alpha}$-Acetyl- $N$-methyl-L-alaninamide and $N^{\alpha}$-Acetyl- $N$-methylglycinamide with the Corresponding 4-21G Values Obtained As Described in the Text

|  | dihedral angles, deg |  |  |  | relative energy, $\mathrm{kcal} / \mathrm{mol}$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Ala |  | Gly |  | Ala |  | Gly |  |
|  | 4-21 G | ECEPP ${ }^{\text {a }}$ | $4-21 G^{\text {b }}$ | ECEPP ${ }^{\text {a }}$ | 4-21G | ECEPP ${ }^{\text {a }}$ | $4-21 \mathrm{G}^{\text {b }}$ | $\overline{\text { ECEPP }}{ }^{\text {a }}$ |
| 1 | $-84.6$ | $-84$ |  | $-83$ | 0.0 | 0.0 | 0.0 | 0.0 |
|  | $73.0$ | $79$ | $70.7$ | $76$ |  |  |  |  |
| II | -165.7 | -154 | -180.0 | -180 | 1.4 | 0.4 | 0.8 | 0.8 |
|  | 167.3 | 153 | 180.0 | 180 |  |  |  |  |
| 1 I | 74.6 | 78 |  |  | 2.6 | 8.8 |  |  |
|  | -62.0 | -64 |  |  |  |  |  |  |
| 1V | -134.2 | -150 | -159.1 | -173 | 3.9 | 0.7 | 4.7 | 1.0 |
|  | 38.1 | 72 | 40.5 | 62 |  |  |  |  |
| V | -91.9 | -74 | $-86.6$ | $-72$ | 4.9 | 1.1 | 4.3 | 1.2 |
|  | $-5.5$ | -45 | -14.1 | --53 |  |  |  |  |
| V1 | 60.8 | 54 |  |  | 6.7 | 2.3 |  |  |
|  | 40.6 | 57 |  |  |  |  |  |  |
| V11 | $\begin{array}{r} -161.7 \\ -55.4 \end{array}$ | $-158$ |  |  | 7.9 | 1.6 |  |  |
|  | $-55.4$ | $-58$ |  |  |  |  |  |  |

[^4]$121.9^{\circ}$ in the $C_{7}{ }^{\text {eq }}$ form and $125.9^{\circ}$ in the $C_{7}{ }^{\text {ax }}$ form, an opening of this angle by $\sim 4^{\circ}$. For the $\mathrm{N} 7-\mathrm{C} 9-\mathrm{C} 12$ and $\mathrm{C} 9-\mathrm{C} 12-\mathrm{N} 17$ angles, the differences are $\sim 2-3^{\circ}$ in the same direction. The opening is probably related to the fact that the $\phi$ and $\psi$ values are smaller in magnitude ( $\phi=74.6^{\circ}$ vs. $\phi=-84.6^{\circ}$ and $\psi=$ $-62.0^{\circ}$ vs. $\psi=73.0^{\circ}$ ) in the $C_{7}^{\text {ax }}$ form than in the $C_{7}^{\text {eq }}$ form. This leads to a more favorable interaction between the $\mathrm{C} 6=\mathrm{O} 5$ and N17-H18 bonds and the side-chain methyl group while, at the same time, it retains optimal hydrogen bonding between $\mathrm{H} 18 \cdots \mathrm{O} 5$. Clearly, the empirical energy calculations do not allow the bond angles to open in the $C_{7}^{\text {ax }}$ form and the nonbonded contact energy is thus not reduced. The result is an energy which may be abnormally high for the $C_{7}^{\text {ax }}$ conformation when calculated by empirical methods in this manner.

Some interesting bond length variations with conformation can be seen from Table II. For example, in $C_{7}{ }^{\text {eq }}$ and $C_{7}^{\text {ax }}$ the $\mathrm{C} 6=\mathrm{O}$ and N17--H18 bonds, which participate in hydrogen bonding are lengthened considerably ( 1.230 (1) $\AA$ and $0.998 \AA$, respectively) compared to the non-hydrogen bonding $\mathrm{C} 12=\mathrm{O} 13$ and $\mathrm{N} 7-\mathrm{H} 8$ (1.225 (6) $\AA$ and $0.995 \AA$, respectively). This effect is to be expected from a charge polarization effect and is in qualitative agreement with previous calculations. ${ }^{1,20,21}$ A second observation that seems to follow the pattern described above is seen for the peptide bond lengths $\mathrm{C} 6-\mathrm{N} 7$ and $\mathrm{C} 12-\mathrm{N} 17$. In this case, the $\omega_{1}$ (C6-N7) bond is shorter for the low-energy structures (I-III) and longer for the higher energy structures (IV-VII). This is the reverse of the $\mathrm{C}=\mathrm{O}$ bond distance trends and in accord with the usual resonance effects. The $\omega_{2}$ (C12-N17) bond shows little variation with change in conformation (except for VII) and is shortest in $C_{7}{ }^{\text {ax }}$.
$C_{5}$ Extended State, The extended form $C_{5}$ (II) represents a particularly flat area of the potential energy surface of Ala. The ab initio energy of the $\mathrm{C}_{5}$ (II) conformation relative to I is higher than that found from ECEPP (see Table III). As pointed out below, this difference may have important implications for the interpretation of the solution data of Ala.

Some additional geometrical properties which are of particular interest can be noted in Table II. For example, the backbone $\mathrm{N} 7-\mathrm{C} 9-\mathrm{C} 12$ bond angle (around the $\mathrm{C}^{\alpha}$ carbon) is smaller $\left(106.4^{\circ}\right)$ in $C_{5}$ than in any other of the calculated structures. The relatively small value for this angle may be an indication of the fact that $\mathrm{Cl} 2=\mathrm{O} 13$ and $\mathrm{N} 7-\mathrm{H} 8$ attract each other in a hydrogen bond type interaction. In this context, it is also interesting to note that the $\mathrm{O} 13-\mathrm{Cl} 2-\mathrm{C} 9$ angle in $C_{5}$ is not significantly smaller than that found in the other conformations. Thus, the smaller bond angle around the $\mathrm{C}^{\alpha}$ may be a result of bond…bond or conjugative effects. The $\mathrm{H} 8 \cdots \mathrm{O} 13$ hydrogen bond length in $C_{5}$ is $2.13 \AA$ and short enough to stabilize this conformation.

The attractive interaction between $\mathrm{N} 7-\mathrm{H} 8$ and $\mathrm{C12}=\mathrm{O} 13$ can also be inferred from a weak correlation between $\phi$ and $\psi$ in the $C_{5}$ conformational region which is apparent from the calculations: During the $4-21 \mathrm{G}$ refinement the optimization followed a path from the starting values of $\phi=-154^{\circ}$ and $\psi=153^{\circ}$ to $\phi=-165^{\circ}$ and $\psi=167^{\circ}$ in which the $\phi$ and $\psi$ values maintained the N7-H8 and $\mathrm{C} 12=013$ bonds in approximately parallel orientation to one a nother.
$\alpha_{\mathrm{R}}$ and $\alpha_{\mathrm{L}}$ States, The positions of minimum energy of the rightand left-handed helix regions found from the 4-21G calculations are shifted in the direction of the bridge structures relative to the equivalent positions found from the empirical ECEPP studies (see Figure 9). In particular, $\alpha_{R}$ is found as an energy trough passing through the bridge ( $\psi \sim 0^{\circ}$ ) region in both the Gly and Ala geometry-optimized calculations (see Table III). This result is a significant finding, since the $\phi, \psi$ values obtained here for these model compounds are obviously closer to $\phi, \psi$ values found in type I (LL or GG) ${ }^{8}$ bends in proteins, than are the values found from empirical calculations. ${ }^{8}$ Examination of Figure 6 shows that in

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Figure 10, $\phi\left(\mathrm{N}-\mathrm{C}^{\alpha}\right)$ and $\psi\left(\mathrm{C}^{\alpha}-\mathrm{C}^{\prime}\right)$ values found from experimental crystal structures of small peptides: (a) ref 25 , (b) ref 26 , (c) ref 27 , (d) ref 28 , (e) ref 29 , (f) ref 30 , (g) ref 31 , (h) ref 24 , (i) ref 32 , (j) ref 33. Structures IV and V are values from Table III.
the $\alpha_{\mathrm{R}}$ conformer the amide hydrogen H 18 lies directly in the region of space occupied by the lone-pair orbital associated with the amide nitrogen N 7 . The $\mathrm{H} 18 \cdots \mathrm{~N} 7$ contact distance in the Ala compound is $2.30 \AA$. This contact distance is quite short, indicating a favorable attractive interaction. In the $\alpha_{L}$ conformer (Figure 7), one also finds the same favorable interaction, but at a distance of $2.42 \AA$. That is, the orientation of the H18 atom relative to the orbital on N7, perpendicular to the plane of the peptide group, is retained in both structures, even though the $\phi$ angle in the $\alpha_{\mathrm{L}}$ conformation cannot become as large in magnitude as it is in $\alpha_{\mathrm{R}}$, probably because of unfavorable interaction between the $\mathrm{C} 5=06$ bond and the side-chain methyl group. Since $\phi$ in $\alpha_{\mathrm{L}}$ remains at $\sim 60^{\circ}$, the $\psi$ value prefers an angle near $\sim 40^{\circ}$, apparently to optimize the overlap between $\mathrm{N} 7 \ldots \mathrm{H} 18$.

Evidence for the importance of the $\mathrm{H} 18 \cdots \mathrm{~N} 7$ interaction is also found from the following observations. First, by draining charge from the N 7 atom one would anticipate that the peptide bond (C6-N7) would be long. Indeed, this bond is $\sim 1.37 \AA$ compared to $1.34-1.35 \AA$ for the $\mathrm{C} 12-\mathrm{N} 17$ bond and $1.35 \AA$ for $\mathrm{C} 6-\mathrm{N} 7$ bonds of structures without this interaction. Furthermore, the $\mathrm{C} 6=\mathrm{O} 5$ bond should become shortened, as should $\mathrm{N} 17-\mathrm{H} 18$ and both are found to be short, in agreement with the overlap effect.

The local geometry effects just described can be used to confirm the hypothesis, originally put forward by Gieren et al. ${ }^{22,23}$ in connection with a crystallographic study of $N, N^{\prime}$-di-tert-butyl-2-[ $N$-(1-phenylethyl)benzamido]malonamide (DTBMA), that a $\phi-\psi$ correlation exists in crystal data of proteins which can be rationalized in terms of a perpendicular $\mathrm{N}-\mathrm{H}, \cdots \mathrm{N}$ hydrogen bond of the type described above. In searching for additional evidence for this important interaction, it can be expected that the concomitant local geometry effects are apparent only with considerable scatter in crystal structures of linear peptides where multiple hydrogen bonding between molecules is found or in larger polymers or proteins, particularly when observed in the helix structure, where hydrogen bonding dominates the conformational energy. However, the trend suggested seems to be found very clearly in structures such as cyclo(L-alanyl-L-alanyl-glycyl-glycyl-L-alanyl-glycyl) monohydrate and cyclo(L-alanyl-L-alanyl-glycyl-L-alanyl-gly-cyl-glycyl) dihydrate ${ }^{24}$. In these structures one finds that as $\phi$

[^6]goes from $\sim-50^{\circ}$ to $\sim-100^{\circ}, \psi$ goes from $\sim-40^{\circ}$ to $\sim 0^{\circ}$. That is, the values of both angles are correlative, in that one does not find any combination such as $\phi \sim-50^{\circ}, \psi \sim 0^{\circ}$ or $\phi \approx-100^{\circ}$, $\psi-40^{\circ}$. The dihydrate also contains an Ala in the $\alpha_{\mathrm{L}}$ conformation ( $\phi=+53.8^{\circ}, \psi=+37.7^{\circ}$ ) indicating the availability of the $\alpha_{\mathrm{L}}$ conforn ation in these cyclic structures. In Figure 10 are plotted the $\phi$ and $\psi$ values found from X-ray structures of a variety of small peptides. No protein data are included, since they have been presented elsewhere. ${ }^{22,23}$ The correlation is quite good throughout the range of angles shown, and remarkably, the low-energy structure IV which also exhibits the $\mathrm{N} 7 \cdots \mathrm{H} 18$ interaction lies on the best fit line shown. The distribution of dihedral angles taking part in protein bends has been presented elsewhere as plots of $\phi$ and $\psi$ (Figure 1 , ref 33 ), and they also show a similar trend, although with considerably more scatter than the data given here. This is not suprising since one would expect long-range as well as steric effects to produce perturbations on these conformational states.

Clearly, the $\mathrm{N}-\mathrm{H} \cdots \mathrm{N}$ interaction defined in this way (such as between N7 and H18 in this study) is a significant factor in peptide structure. It appears to be an important contribution to the formation and stability of bend conformations. Long-range interactions undoubtedly also play a role in stabilizing bend structures in proteins, as do the constraints imposed by cyclization of medium-sized polypeptides. However, the dominance of this interesting local interaction in dipeptide sequences must have significant implications in the mechanism of folding of proteins and in the overall stability of native protein structures.

The $\alpha_{\mathrm{R}}$ local geometry is somewhat perturbed from the average of the other conformations with the bond angle around the $\mathrm{C}^{\alpha}$ ( $\mathrm{N} 7-\mathrm{C} 9-\mathrm{C} 12$ ) expanded to $114.0^{\circ}\left(111.1^{\circ}\right.$ in $\left.\alpha_{\mathrm{L}}\right)$ in the Ala and $115.2^{\circ}$ in the Gly molecule. ${ }^{1}$ The side-chain methyl group in the Ala molecule seems also to have moved toward the neighboring carbonyl group in the $\alpha_{\mathrm{R}}$ structure as shown by the small C11-C9-C12 angle (i.e., $108.3^{\circ}$ ). This angle is smaller in $\alpha_{\mathrm{R}}$ than in any other conformation.

It is also most compelling to note that, in the crystal structure of DTBMA, ${ }^{22.23}$ the $\mathrm{H}, \cdots \mathrm{N}$ contact distance is $2.27 \AA$ in excellent agreement with the $2.30 \AA$ found here. The peptide bond length of the nitrogen involved in this interaction is also longer ( $1.37 \AA$ ) than those peptide bonds which are not involved.

Comparison with Solution Studies, Blocked single residue molecules have been the subject of several experimental ${ }^{36}$ studies, and several interpretations of the spectra have been given. ${ }^{36}$ The spectral analysis has usually been based on the presence of two conformation in solution: the $C_{5}$ conformation with $(\phi, \psi) \simeq$ $\left(-170^{\circ}, 170^{\circ}\right)$ and the $C_{7}{ }^{\text {eq }}$ conformation with $(\phi, \psi) \simeq\left(-80^{\circ}\right.$, $\left.80^{\circ}\right)$. A third conformation, called $\gamma$ with $(\phi, \psi) \simeq\left(-60^{\circ}, 140^{\circ}\right)$ has also been proposed. ${ }^{37}$ Recently, ${ }^{38}$ it was noted that in chloroform solution, no evidence for the strongly hydrogen-bonded $C_{7}{ }^{\text {eq }}$ could be found, even though absorption bands attributable to $C_{5}$ were observed.

From the calculated ab initio energies (Table III), one would expect $C_{7}{ }^{\text {eq }}$ and $C_{5}$ to predominate in solution. However, the actual populations of the $C_{7}^{\text {eq }}$ and $C_{s}$ states in solution are also affected by unknown entropic terms. If the ECEPP potential energy surface (Figure 9) gives a qualitatively reasonable description of

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Figure 11, Geometry around the peptide bond. The ab initio values listed are the minimum and maximum values found from seven conformers of Ala optimized. Experimental values are given in parentheses.
the Ala system, then it can be inferred that there is a larger number of low-lying states in the vicinity of $C_{5}$ than of $C_{7}{ }^{\text {eq }}$, whereas very few low-lying energy states should be available in the vicinity of $C_{7}{ }^{\text {ax }}$, which is also one of the relatively stable forms of this system. The higher ab initio energy of $C_{5}$ in Ala as compared to $C_{5}$ in Gly relative to $C_{7}{ }^{\text {eq }}$ seems to be reasonable because the $C_{5} / C_{7}{ }^{\text {eq }}$ ratios inferred from the solution spectra ${ }^{38}$ are $\sim 0.4$ for Ala and 0.7 for Gly.
Geometry of the Peptide Unit, The variation in geometry found here as a function of conformation leads one to approach cautiously any comparison of calculated bond lengths and angles to those found experimentally. In making such comparisions it must also be reiterated that ab initio calculations do not reproduce experimental structures exactly but that there are characteristic constant differences for parameters of a given type. ${ }^{1,15,18 \mathrm{c}}$ A summary of such differences for $4-21 \mathrm{G}$ structures is given in ref 1 and 18 c .

For a comparison of calculated and experimental geometric parameters around the peptide bond, we will use here a recent summary of crystal data of some 20 linear peptide derivatives with trans-peptide units. ${ }^{39}$ This summary refers to different amino acids in different conformational states which are in addition also affected by differences in the crystal environment. Variations in local geometry with conformational arrangement are, therefore, averaged out in the experimental data.

The calculated peptide geometry shown in Figure 11 is in remarkably good agreement with the experimental average. The $a b$ initio values bracket the experimental results in most cases, the major exception being the $\mathrm{C}^{\prime}-\mathrm{N}$ (peptide) bond length where the calculated values are always longer than the crystal average. Interestingly, long $\mathrm{C}^{\prime}-\mathrm{N}$ bonds are also found from gas-phase electron diffraction studies of formamide ( $r_{\mathrm{CN}}=1.368 \AA$ ), ${ }^{40}$ methylformamide $\left(r_{\mathrm{CN}}=1.366 \AA\right),{ }^{41}$ acetamide $\left(r_{\mathrm{CN}}=1.380 \AA\right)$, ${ }^{42}$ and $N$-methylacetamide $\left(r_{\mathrm{CN}}=1.386 \AA\right) .{ }^{43} \quad$ The $\mathrm{C}=\mathrm{O}$ bonds in the same series are all relatively short $\left(r_{\mathrm{CO}}=1.212 \AA,{ }^{40} 1.219\right.$ $\AA,{ }^{41} 1.220 \AA,{ }^{42}$ and $1.225 \AA,{ }^{43}$ respectively). It seems to be a general trend that $\mathrm{C}=\mathrm{O}$ bond distances are always longer in the solid state than in the gas phase, whereas $\mathrm{C}-\mathrm{N}$ bonds in the crystalline state are shorter than in the vapor phase. Since the gas-phase results are also consistent with the ab initio calculations, one might suggest that these parameters are affected by their environment in the crystal.

Peptide Bond ( $\omega_{1}$ and $\omega_{2}$ ), The two peptide groups (see Table II) are nonplanar in all the Ala conformers, but deviations from planarity can be nearly negligible in some structures as, for example, in II. As discussed previously, the backbone conformation

[^8]plays an important role in the degree of nonplanarity of $\omega_{1}$. Interestingly, the ECEPP calculations do not show significant deviations from planarity in $\omega$ resulting from different low-energy backbone conformations ${ }^{34}$ even though upon N -methylation large deviations are found. ${ }^{35}$

Comparison to Other ab Initio Calculations, A comparison of the results presented here with other ab initio calculations of $\mathrm{Ala}^{3}$ shows one very obvious difference. That is, in the $\phi-\psi$ map prepared without geometry optimization on the STO-3G level, ${ }^{3}$ the $C_{7}{ }^{\text {ax }}$ conformation does not appear as a region of low energy. We believe that this difference is an artifact of refinement related to the special nature of $C_{7}{ }^{\text {ax }}$ as a narrow cavity in the Ala potential energy surface.

As indicated in Figure 9, $C_{7}{ }^{\text {ax }}$ is a narrow potential energy minimum in the ECEPP calculation. Even though our torsional angles are somewhat uncertain for the reasons pointed out above and even though the space around this conformation has not been examined here on the $4-21 \mathrm{G}$ level, it is not impossible that this conformer was missed in the $20^{\circ}$ grid used to obtain the $\phi-\psi$ map. ${ }^{3}$ In addition, the rather large geometry change (see above) is required for this structure to be optimized. Thus, conformational problems inherent in the use of rigid geometry appear to be present whether empirical methods or the more rigorous ab initio calculations are used. The other conformational regions of low energy found in the STO-3G $\phi-\psi$ map $^{3}$ are relatively close to the characteristic areas considered here.

## Conclusion

The geometry-optimized conformations described here clearly show the structural changes that must be considered in future conformational studies of peptides. Some geometry changes are very significant, such as that found for the $C_{7}{ }^{\text {ax }}$ conformer, while others are more subtle. We consider the perpendicular peptide interaction between the nitrogen N 7 and second amide hydrogen H18 to be a most significant and important factor for the understanding of dipeptide conformation. The calculated details of local geometry make it very likely that this interaction is responsible for the observed correlation (see Figure 10) between the $\phi$ and $\psi$ torsional angles in the $\alpha$-helix conformational region and must play an important role in the formation and stability of protein structure. Its role in bend structures is being examined, and its incorporation into empirical calculations should lead to improved conformational predictions of larger polypeptides.

Acknowledgment. We thank Dr. F. R. Helm, Dr. B. Ashmore, Ron Fowler, C. Stromberg, and J. Swayze of Computing Services for their assistance and the University of Arkansas Computing Center for executing the calculations. This research was supported in part by NSF Grants CHE-7920790 and ISP-8011447. C.V.A. is grateful for a NATO Research Fellowship. L.S. acknowledges helpful discussions with Prof. P. Pulay (Arkansas). The technical assistance of Ms. Monica L. Mabie is gratefully acknowledged.

Registry No, $N^{\alpha}$-Acetyl- $N$-methylalaninamide, 19701-83-8.

# Gas-Phase Derivatization for Determination of the Structures of $\mathrm{C}_{3} \mathrm{H}_{5}{ }^{+}$Ions 

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#### Abstract

The structure of the $\mathrm{C}_{3} \mathrm{H}_{5}{ }^{+}$ions formed via halide atom loss from ionized allyl, cyclopropyl, and 1- and 2-propenyl halides was investigated by derivatizing the ions with neutral benzene and substituted benzenes to produce gas-phase adduct ions. The structures of pressure stabilized adduct ions were directly determined by obtaining their collision-induced decomposition spectra and comparing them with the CID spectra of model ions. Two $\mathrm{C}_{3} \mathrm{H}_{5}{ }^{+}$ion structures were found to be stable, the allyl cation from cyclopropyl and allyl halides and the 2-propenyl cation from 2-propenyl halides. The $\mathrm{C}_{3} \mathrm{H}_{5}{ }^{+}$ions from 1-propenyl halides exist as a $2: 1$ mixture of the two ion structures. The mechanism of reaction was shown to be electrophilic attack on the ring $\pi$ system to produce a Wheland intermediate or other structure with the proton relocated. The reaction pathway was confirmed by Fourier transform mass spectrometry, and rate constants for the derivatization reaction were determined by using pulsed ICR.


Vinyl and allyl cations have been the subjects of considerable study both in solution ${ }^{1}$ and in the gas phase. In solution, the ions are important as reactive intermediates in synthesis ${ }^{2}$ and as models for rearrangement reactions. ${ }^{3}$ However, direct observation has been limited to studies of appropriately substituted ions or those formed via loss of "super" leaving groups. ${ }^{4}$ Moreover, their

[^9]reactivity in solution may largely reflect the nature of the solvent system rather than the properties of the ions themselves.

The simplest vinyl-allyl system, namely, the $\mathrm{C}_{3} \mathrm{H}_{5}{ }^{+}$isomers and in particular the allyl cation, has been the subject of a number of recent studies in the gas phase. The vertical and adiabatic ionization potential (IP) of the allyl radical as well as the appearance potential (AP) of allyl cations were measured. ${ }^{5}$ These measurements lead to a heat-of-formation of $226 \mathrm{kcal} / \mathrm{mol}$, which has been independently verified. ${ }^{6}$ The heats of formation of numerous $\mathrm{C}_{3} \mathrm{H}_{5}{ }^{+}$fragment ions from simple hydrocarbons were determined to be identical with the value for allyl cation. ${ }^{7}$

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